

# Entanglement and spontaneous parametric down conversion

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$$\hat{\mathcal{H}} = \sum_{j=p,s,i} \hbar\omega_j \left( \hat{n}_j + \frac{1}{2} \right) + \hbar g \left( \hat{a}_s^\dagger \hat{a}_i^\dagger \hat{a}_p + cc \right) \quad (1)$$

$\hat{n}_j = \hat{a}_j^\dagger \hat{a}_j$  photon number operator

$g$  coupling constant, including the nonlinear susceptibility

$$\hat{\mathcal{H}}(t) = \frac{\epsilon_0}{2} \int_V d^3r \chi^{(2)} \hat{E}_p(r, t) \hat{E}_s(r, t) \hat{E}_i(r, t)$$

$$\hat{\mathbf{E}}^{(+)} = i \sum_{\ell} \sqrt{\frac{\hbar \omega_{\ell}}{2 \epsilon_0 L^3}} \vec{\epsilon}_{\ell} \hat{a}_{\ell} e^{i(\mathbf{k}_{\ell} \cdot \mathbf{r} - \omega_{\ell} t)}$$

$$\begin{aligned} \hat{\mathcal{H}}_{int}(t) \propto & \frac{1}{L^3} \epsilon_0 \sum_{\ell} \sum_{\ell'} \mathcal{A}_p \vec{\epsilon}_p \chi^{(2)}(\omega_p, \omega_{\ell}, \omega_{\ell'}) \vec{\epsilon}_{\ell} \vec{\epsilon}_{\ell'} \int_V d^3r \hat{a}_{\ell} \hat{a}_{\ell'} \\ & \times e^{i(\mathbf{k}_p - \mathbf{k}_{\ell} - \mathbf{k}_{\ell'}) \cdot \mathbf{r}} e^{i(\omega_{\ell} + \omega_{\ell'} - \omega_p)t} + H.c. \end{aligned}$$

$$i\hbar \frac{d|\psi\rangle}{dt} = \hat{\mathcal{H}}_{int}(t)|\psi\rangle$$

$$|\psi(t)\rangle = \exp\left(\frac{1}{i\hbar} \int_0^t \hat{\mathcal{H}}_{int}(t') dt'\right) |\psi(0)\rangle$$

with  $|\psi(0)\rangle = |0\rangle_\ell |0\rangle_{\ell'}$

# Spontaneous parametric down conversion photon pair state

$$|\psi(t)\rangle \simeq \left( 1 + \frac{1}{i\hbar} \int_0^t \hat{\mathcal{H}}_{int}(t') dt' \right) |\psi(0)\rangle$$

$$|\psi(t)\rangle \simeq |\psi(0)\rangle + C \sum_{\ell} \sum_{\ell'} \Phi(\Delta k_{\ell, \ell'}, \omega_{\ell}, \omega_{\ell'}, \omega_p) \delta(\omega_{\ell} + \omega_{\ell'} - \omega_p) |1\rangle_{\ell} |1\rangle_{\ell'}$$

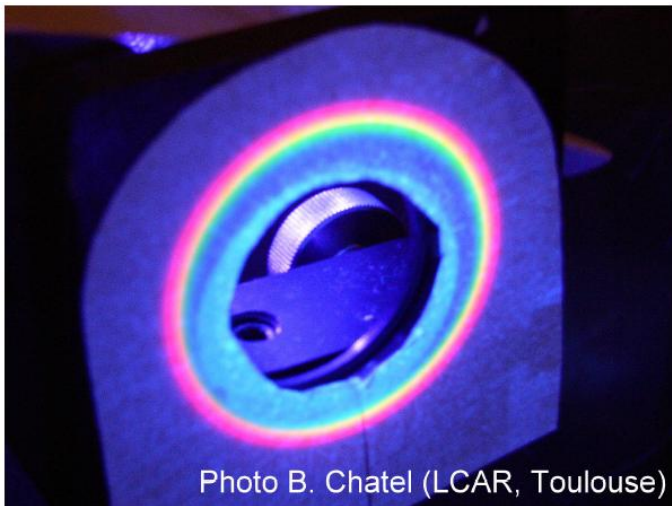
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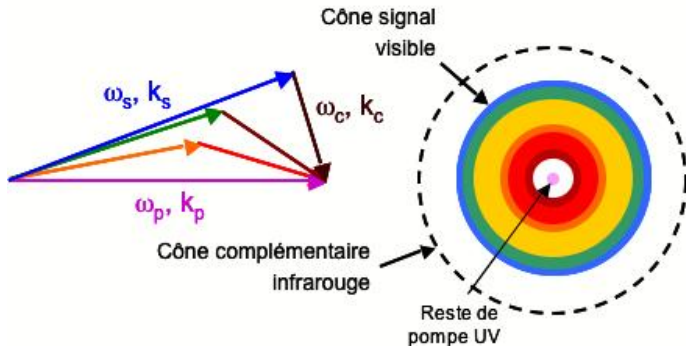
entangled state

# Spontaneous parametric down conversion



$$\lambda_P = 400\text{nm}$$

# Spontaneous parametric down conversion





# Franson Bell experiment

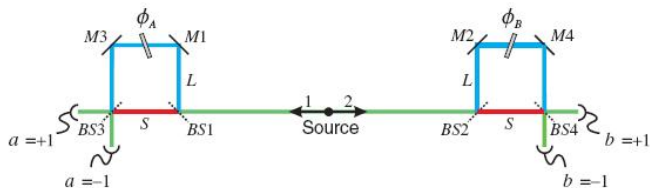


FIG. 1 (color online). Generic setup of the Franson Bell experiment.

# Franson Bell experiment

source can be

- atomic cascade
  - first excited level : long lifetime  $\tau_1$
  - intermediate level:  $\tau_2 \ll \tau_1$

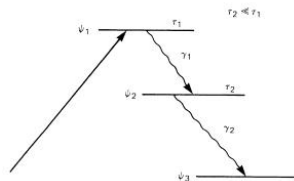


FIG. 1. Three-level atomic system with a relatively long lifetime  $\tau_1$  for the initial state and a much shorter lifetime  $\tau_2$  for the intermediate state.

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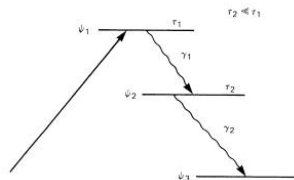


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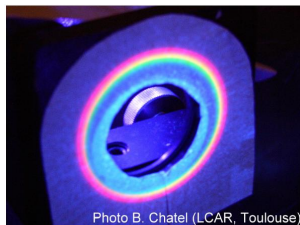


Photo B. Chatel (LCAR, Toulouse)

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- emission time
- time difference between the two photons

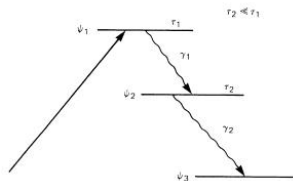


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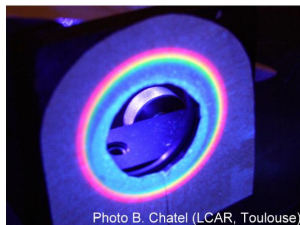


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- spontaneous conversion
  - $\tau_1 \longleftrightarrow \mathcal{T}_{\text{pump}}^c$
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in both case,

- emission time **unknown**
- time difference between the two photons **precisely known**

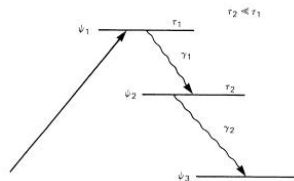
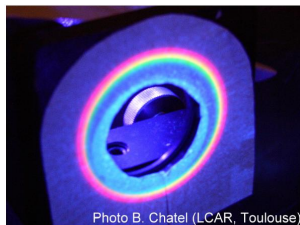


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# Franson Bell experiment: interferences

Indistinguishability between

$$|S\rangle_A \otimes |S\rangle_B \text{ and } |L\rangle_A \otimes |L\rangle_B$$

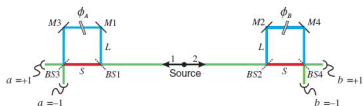


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## Feynman and interferences in quantum physics

Indistinguishable paths  $\longleftrightarrow$  interference

- A given path ( $|s\rangle_A \otimes |s\rangle_B$ )  $\longrightarrow$  given result with a probability  $|\alpha_A|^2$
- Another path ( $|l\rangle_A \otimes |l\rangle_B$ )  $\longrightarrow$  same result with a probability  $|\alpha_B|^2$
- If the paths are distinguishable, the probability of the result is  $|\alpha_A|^2 + |\alpha_B|^2$  ( incoherent sum)
- If the paths are indistinguishable, the probability of the result is  $|\alpha_A + \alpha_B|^2$  (coherent sum)  $\longrightarrow$  **the relative phase is important**

Coincidence rate

$$R_c = \frac{1}{4} R_0 \cos^2 \left( \frac{\Delta E \Delta T / \hbar + \Phi_A + \Phi_B}{2} \right)$$

$\Delta T$  transit time difference between short and long arm of the interferometer ( $\Delta L$ )

| physical process | $\Delta T$  | $\Delta E$                  |
|------------------|---|-----------------------------|
| atomic cascade   | $\tau_2 \ll \Delta T \ll \tau_1$                          | $E_3 - E_1$                 |
| SPDC             | $T_{\text{pump}}^c \gg \Delta T \gg T_{\text{photons}}^c$ | $\hbar\omega_{\text{pump}}$ |

# Synchronization issue

Possible only in the case of spontaneous down conversion in a nonlinear crystal

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<sup>1</sup>Tanzilli,  $\lambda=710$  nm

<sup>2</sup>Tanzilli,  $\lambda=657$  nm



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- OK with CW laser  $\mathcal{L}_{\text{pompe}}^c \sim 300 \text{ m}^1$  ( $\Delta L \sim \text{qq cm}$ )
- impossible with a pulsed laser : 400 ps  $\rightarrow \mathcal{L}_{\text{pompe}}^c \sim 140 \mu\text{m}^2$   
( $\Delta L \sim \mu\text{m}$ )

---

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# Synchronization issue

With a pulsed pump

$|S\rangle_A \otimes |S\rangle_B$  and  $|L\rangle_A \otimes |L\rangle_B$   
are no longer indistinguishable

because  $\mathcal{T}_{\text{pulsedpump}}^c \ll \Delta T$

With a pulsed pump

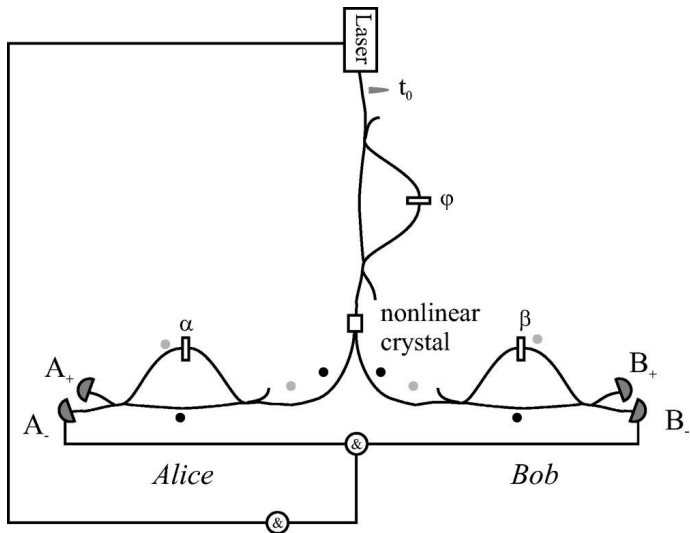
$|S\rangle_A \otimes |S\rangle_B$  and  $|L\rangle_A \otimes |L\rangle_B$   
are no longer indistinguishable

because  $T_{\text{pulsedpump}}^c \ll \Delta T$

Restores undetermination of the emission time

# Synchronization issue : time-bin entanglement

## Pump preparation





because

$$\Delta L \gg \mathcal{L}_c(\text{pulsed pump})$$

The pump photon has passed

- **either** through the short path
- **or** through the long path

States that can be produced

$$\frac{1}{\sqrt{2}} \left( |S\rangle_P |S\rangle_P + e^{i\phi} |L\rangle_P |L\rangle_P \right)$$
$$\frac{1}{\sqrt{2}} \left( |S\rangle_P |L\rangle_P + e^{i\phi} |L\rangle_P |S\rangle_P \right)$$

# Time-bin entanglement

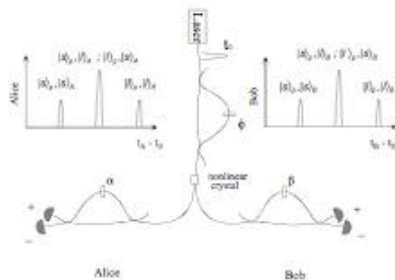


FIG. 1. Schematics of quantum key distribution using energy-time Bell states.

# Two independent bases

- Time basis

$|SS\rangle_A \otimes |SS\rangle_B$  (left satellite time-bin)

$|LL\rangle_A \otimes |LL\rangle_B$  (right satellite time-bin)

Time correlation

## Time basis

→ "Particle like behavior"

- Energy basis

$|SL\rangle_A \otimes |SL\rangle_B$  or  $|LS\rangle_A \otimes |LS\rangle_B$  indistinguishable

Interferences

$P_{i,j} = \frac{1}{4} (1 + ij \cos(\alpha + \beta - \phi))$  where  $i, j = \pm 1$  (detector labels)

## Energy basis

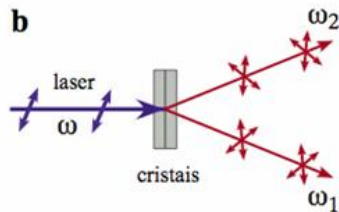
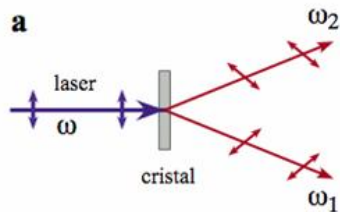
→ "Wave like behavior"

# Polarisation entanglement

$$\frac{1}{\sqrt{2}} \left( |H\rangle_A |H\rangle_B + e^{i\phi} |V\rangle_A |V\rangle_B \right)$$

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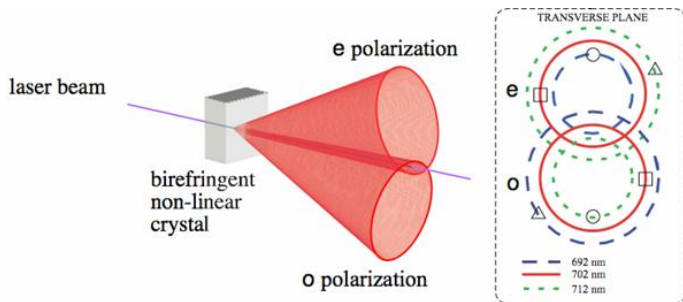
Type I phase matching

# Polarisation entanglement

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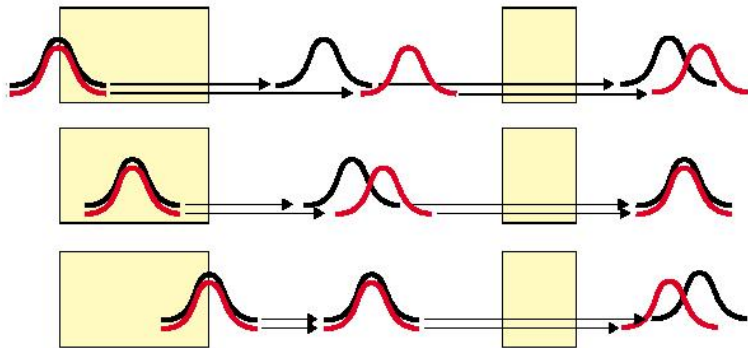
# Polarisation entanglement

$$\frac{1}{\sqrt{2}} \left( |H\rangle_A |V\rangle_B + e^{i\phi} |V\rangle_A |H\rangle_B \right)$$



## Type II phase matching

# Temporal walk-off and compensation





# Temporal walk-off and compensation

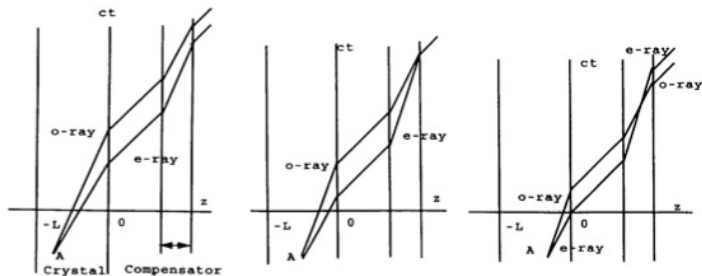
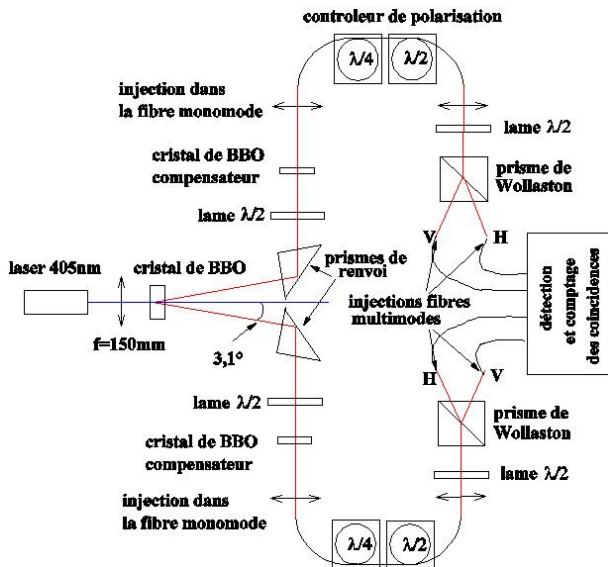


FIG. 6. These diagrams illustrate how the compensator effects pairs that are created at point A near the input, at the center, and near the output of the crystal.

# Polarisation entanglement



# Polarisation entanglement

